



# Mass Spectrum dependence of Higgs-mediated $\mu$ -e Transition in the MSSM

Based on [Phys.Lett.B694:380-385,2011](#)  
[PhysRevD.84.033011](#)

Masaki J. S. YANG  
(Univ. of Tokyo)

w/ Junji Hisano, Shohei Sugiyama, Masato Yamanaka



## Mass Spectrum dependence of Higgs-mediated $\mu$ -e Transition in the MSSM

- CLFV in the MSSM
- The Higgs mediated LFV
- Conclusion

# cLFV in the New Physics

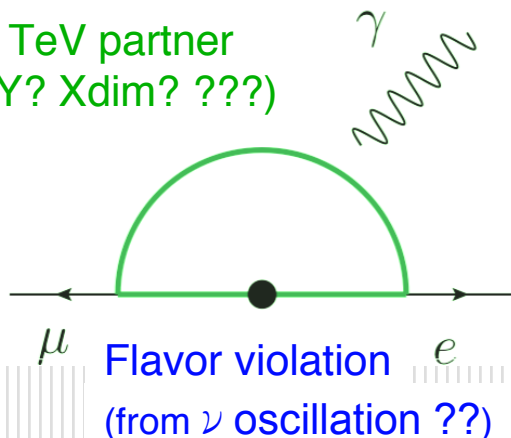
In the SM, there is **no LFV**. However,

Discovery of neutrino oscillation (SK, KamLand, SNO, K2K, etc,,)

TeV physics requires partner of SM particles (SUSY? Xdim? LH???)

$\mu \rightarrow e \gamma$  decay

some TeV partner  
(SUSY? Xdim? ???)



$$\text{BR}(\mu \rightarrow e \gamma) \sim \frac{\alpha}{4\pi} \frac{m_W^4}{(\Lambda \sim \text{TeV})^4} \Delta_{\mu e}^2 F_{\text{NP}}$$

Perhaps, **detectable**?

# cLFV in the New Physics

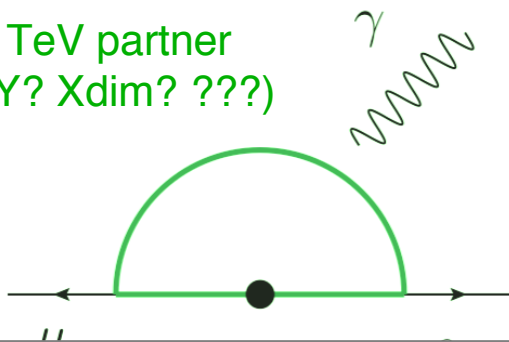
In the SM, there is **no LFV**. However,

Discovery of neutrino oscillation (SK, KamLand, SNO, K2K, etc,,)

TeV physics requires partner of SM particles (SUSY? Xdim? LH???)

$\mu \rightarrow e \gamma$  decay

some TeV partner  
(SUSY? Xdim? ???)



$$\text{BR}(\mu \rightarrow e \gamma) \sim \frac{\alpha}{4\pi} \frac{m_W^4}{(\Lambda \sim \text{TeV})^4} \Delta_{\mu e}^2 F_{\text{NP}}$$

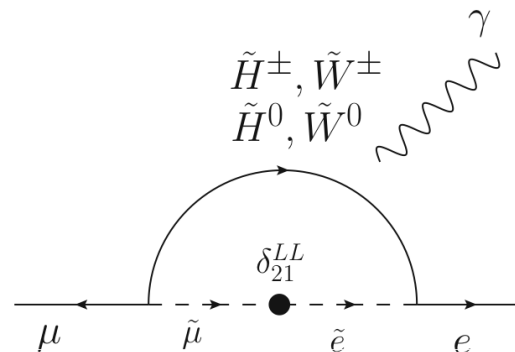
cLFV process is clean signature of beyond the SM!

# cLFV in the MSSM

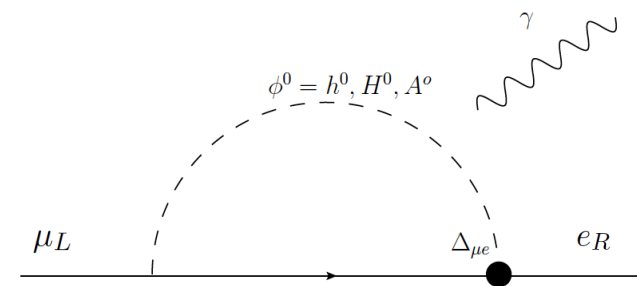
- In the SUSY models, these cLFV are induced by
  - Ordinary slepton / gaugino exchange
  - Higgs exchange (Babu&Kolda, 2000, 2002)  
(Masiero et al, Brignole & Rossi, Ellis et al, etc,,)

$$\mu \rightarrow e \gamma$$

Gaugino LFV

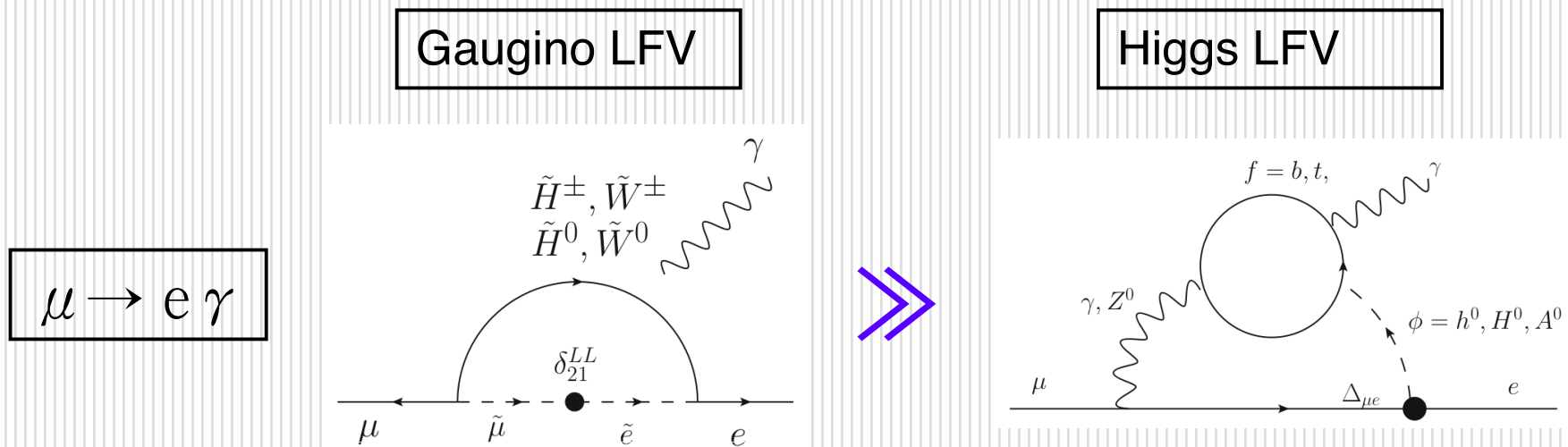


Higgs LFV



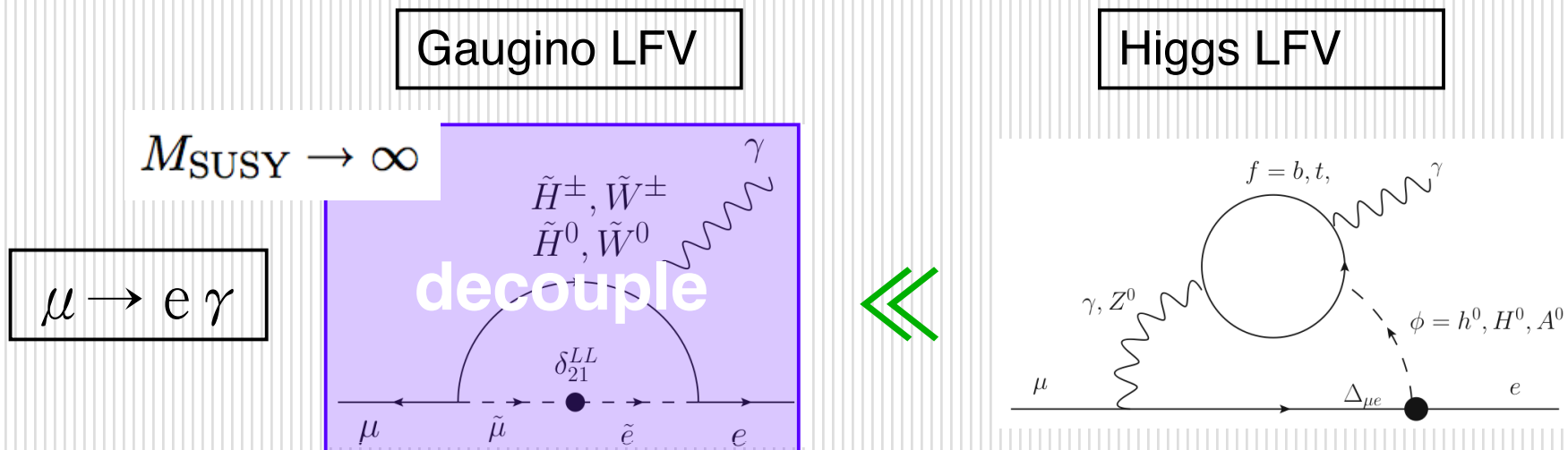
# cLFV in the MSSM

- In the SUSY models, these cLFV are induced by
  - Ordinary slepton / gaugino exchange
  - Higgs exchange (Babu&Kolda, 2000, 2002)  
(Masiero et al, Brignole & Rossi, Ellis et al, etc,,,) )
- Especially  $\mu \rightarrow e$ , **two-loop effects** can be **dominant**  
(Paradisi, 2005, 2006)



# cLFV in the MSSM

- We calculate this two-loop (Barr-Zee) dgms. and  $\text{BR}(\mu\text{Al} \rightarrow e\text{Al}) / \text{BR}(\mu \rightarrow e\gamma)$  in some mass parameters.
  - It is considered that the Higgs effect becomes dominant when SUSY particles are decouple (non decoupling effect).

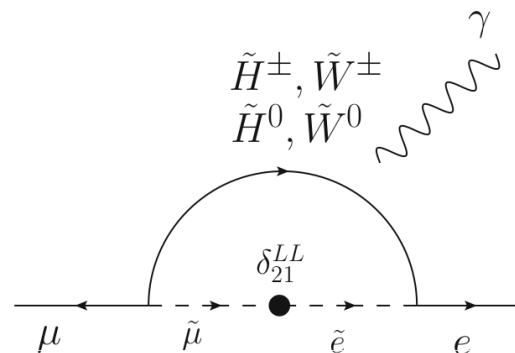


# cLFV in the MSSM

- We calculate this two-loop (Barr-Zee) dgms. and  $\text{BR}(\mu \text{Al} \rightarrow e \text{Al}) / \text{BR}(\mu \rightarrow e \gamma)$  in some mass parameters.
  - It is considered that the Higgs effect becomes dominant when SUSY particles are decouple (non decoupling effect).
- However, if there are destructive interference, it is necessary to consider the Higgs LFV effect.

$$\mu \rightarrow e \gamma$$

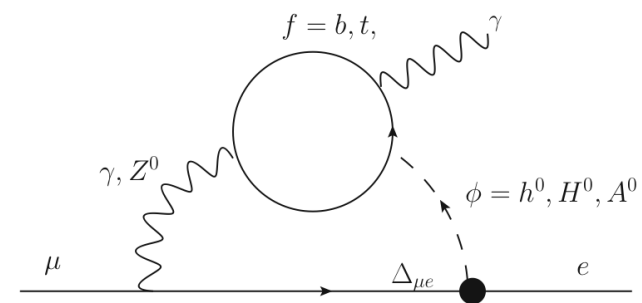
Gaugino LFV



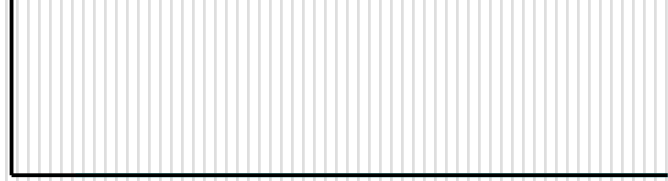
!



Higgs LFV







# The Higgs mediated LFV

# Higgs mediated FCNC

- In the SUSY model, Superpotential must be **holomorphic**.

$$W_{\text{MSSM}} = Y_u U^c Q H_2 - Y_d D^c Q H_1 - Y_e E^c L H_1 + \mu H_1 H_2$$

SUSY  $\Rightarrow \bar{e}_{Ri} Y_{ij}^e e_{Lj} H_d^0 \Rightarrow$  1 Higgs for 1 fermion

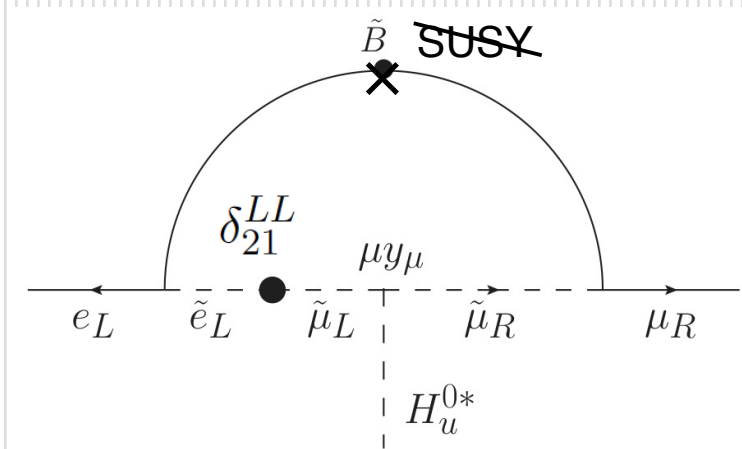
$\rightarrow$  **no FCNC** in the Higgs sector

But, ~~SUSY~~  $\Rightarrow \bar{e}_{Ri} \Delta_{ij} e_{Lj} H_u^{0*} \Rightarrow$  2 Higgs for 1 fermion

$\rightarrow$  **off diag Yukawa & FCNC occurs radiatively**

# Higgs mediated LFV in the MSSM

- LFV Yukawa int. is induced radiatively.



Babu, Kolda (2000,2002)

Diagram illustrating the effective LFV Yukawa interaction. The muon line ( $\mu_L, \mu_R$ ) is connected to the electron line ( $e_L, e_R$ ) via a Higgs boson ( $H_u^{0*}$ ) exchange. The vertex is labeled  $y_i \Delta_{ij}^L$ .

$$\bar{e}_{Ri} y_i \Delta_{ij}^L e_{Lj} H_u^{0*}$$

Assuming all the SUSY mass are degenerate :

$$\Delta_{21}^L \sim \frac{\alpha_2}{24\pi} \delta_{21}^{LL} \quad \text{independent to } M_{\text{SUSY}} \text{ in first order}$$

⇒ Although SUSY particles mass ( $M_{\text{SUSY}}$ ) is large, this effects **remains** (non decoupling effects).

# Set up of the study

## The Non universal Higgs Model (NUHM)

<div><math>m_{H_u}</math></div> <div><math>m_{H_d}</math></div>	quark	$q$	1/2	spin	0	squark	$\tilde{q}$	$m_0$
	lepton	$l$	1/2		0	slepton	$\tilde{l}$	
	gauge boson	$A_\mu$	1		1/2	gaugino	$\lambda$	$m_{1/2}$
	Higgs boson	$H_u$	0	$\mu$	1/2	Higgsino	$\tilde{H}_u$	
		$H_d$		$m_{A^0}$	By EWSB condition;	$\tilde{H}_d$		

And we assume either left or right handed flavor mixing;

$$(m_{\tilde{e}_L}^2)_{ij} = m_{\tilde{e}_L}^2 (1 + \delta_{ij}^{LL}) \quad (m_{\tilde{e}_R}^2)_{ij} = m_{\tilde{e}_R}^2 (1 + \delta_{ij}^{RR})$$

# Set up of the study

## The Non universal Higgs Model (NUHM)

	quark	$q$	1/2	spin	0	squark	$\tilde{q}$	$m_0$
	lepton	$l$	1/2		0	slepton	$\tilde{l}$	
	gauge boson	$A_\mu$	1		1/2	gaugino	$\lambda$	$m_{1/2}$
$m_{H_u}$	Higgs boson	$H_u$	0	$\mu$	1/2	Higgsino	$\tilde{H}_u$	By EWSB condition;
$m_{H_d}$		$H_d$		$m_{A^0}$			$\tilde{H}_d$	

And we assume either left or right

Free parameter

$$(m_{\tilde{e}_L}^2)_{ij} = m_{\tilde{e}_L}^2 (1 + \delta_{ij}^{LL})$$

$$(m_{\tilde{e}_R}^2)_{ij} = m_{\tilde{e}_R}^2 (1 + \delta_{ij}^{RR})$$

# Set up of the study

## The Non universal Higgs Model (NUHM)

<div><math>m_{H_u}</math></div> <div><math>m_{H_d}</math></div>	quark	$q$	1/2	spin	0	squark	$\tilde{q}$	$m_0$
	lepton	$l$	1/2		0	slepton	$\tilde{l}$	
	gauge boson	$A_\mu$	1		1/2	gaugino	$\lambda$	$m_{1/2}$
	Higgs boson	$H_u$	0	$\mu$	1/2	Higgsino	$\tilde{H}_u$	
		$H_d$		$m_{A^0}$	By EWSB condition;	$\tilde{H}_d$		

And we assume either left or right handed flavor mixing;

$$(m_{\tilde{e}_L}^2)_{ij} = m_{\tilde{e}_L}^2 (1 + \delta_{ij}^{LL})$$

SUSY seesaw

$$(m_{\tilde{e}_R}^2)_{ij} = m_{\tilde{e}_R}^2 (1 + \delta_{ij}^{RR})$$

SUSY SU(5) GUT

# Set up of the study

## The Non universal Higgs Model (NUHM)

quark	$q$	1/2	spin	0	squark	$\tilde{q}$	$m_0$
lepton	$l$	1/2		0	slepton	$\tilde{l}$	$m_{1/2}$
<div>We discussed <math>\mu</math>-e transition from the gaugino / Higgs effects.</div>							
$m_{H_u}$	$H_d$	$m_{A^0}$	condition;	$H_d$			
$m_{H_d}$							

And we assume either left or right handed flavor mixing;

$$(m_{\tilde{e}_L}^2)_{ij} = m_{\tilde{e}_L}^2 (1 + \delta_{ij}^{LL})$$

SUSY seesaw

$$(m_{\tilde{e}_R}^2)_{ij} = m_{\tilde{e}_R}^2 (1 + \delta_{ij}^{RR})$$

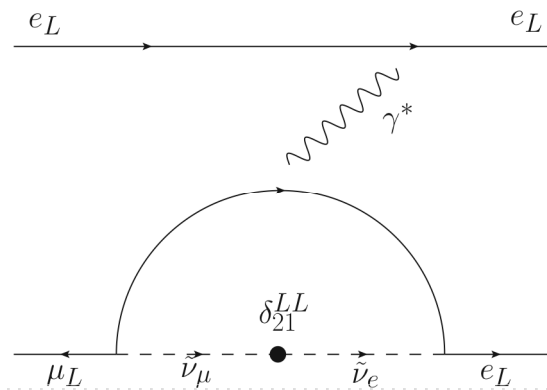
SUSY SU(5) GUT

# Correlation between $\mu$ - e conversion

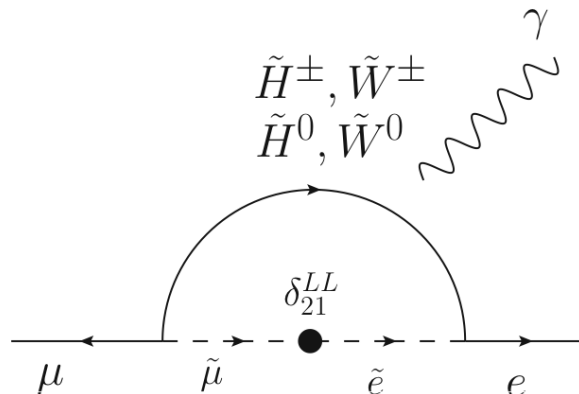
- Evaluate  $\text{BR}(\mu A1 \rightarrow e A1) / \text{BR}(\mu \rightarrow e \gamma)$  in some mass parameters.

## Gaugino LFV

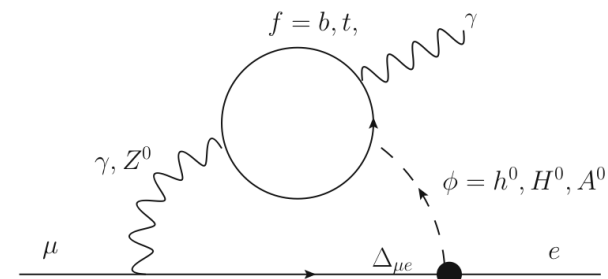
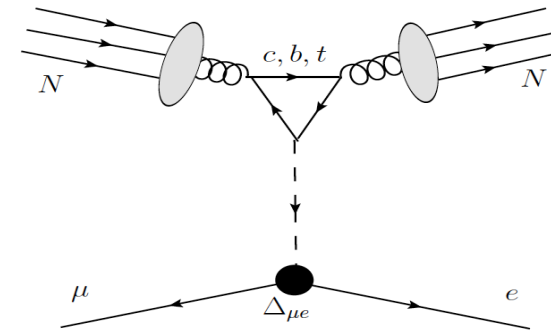
$\mu A1 \rightarrow e A1$



$\mu \rightarrow e \gamma$



## Higgs LFV





# Correlation between $\mu$ - $e$ conversion

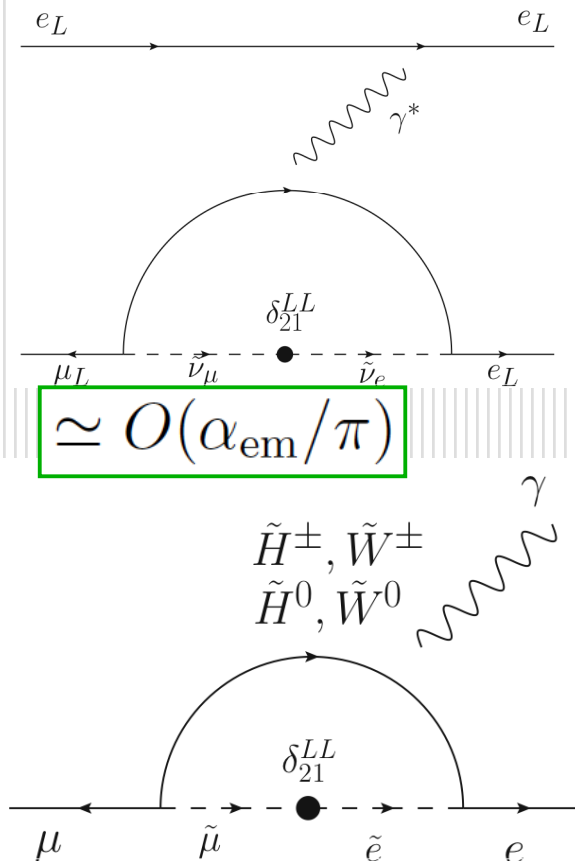
- Evaluate  $\text{BR}(\mu\text{Al} \rightarrow e\text{Al}) / \text{BR}(\mu \rightarrow e\gamma)$  in some mass parameters.

## Gaugino LFV

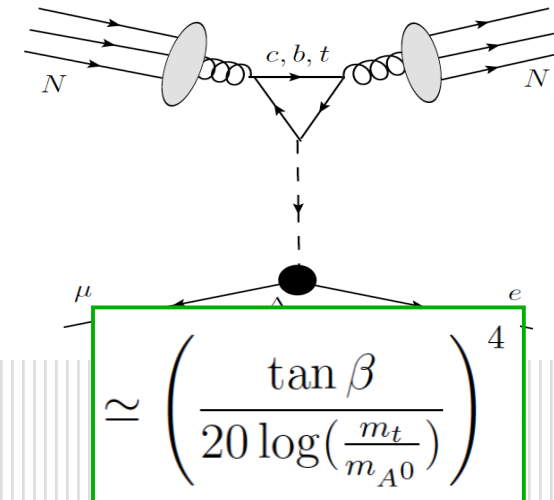
$$\mu\text{Al} \rightarrow e\text{Al}$$

$$\frac{\text{BR}(\mu\text{Al} \rightarrow e\text{Al})}{\text{BR}(\mu \rightarrow e\gamma)}$$

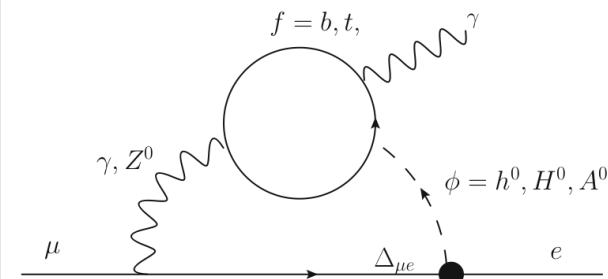
$$\mu \rightarrow e\gamma$$



## Higgs LFV

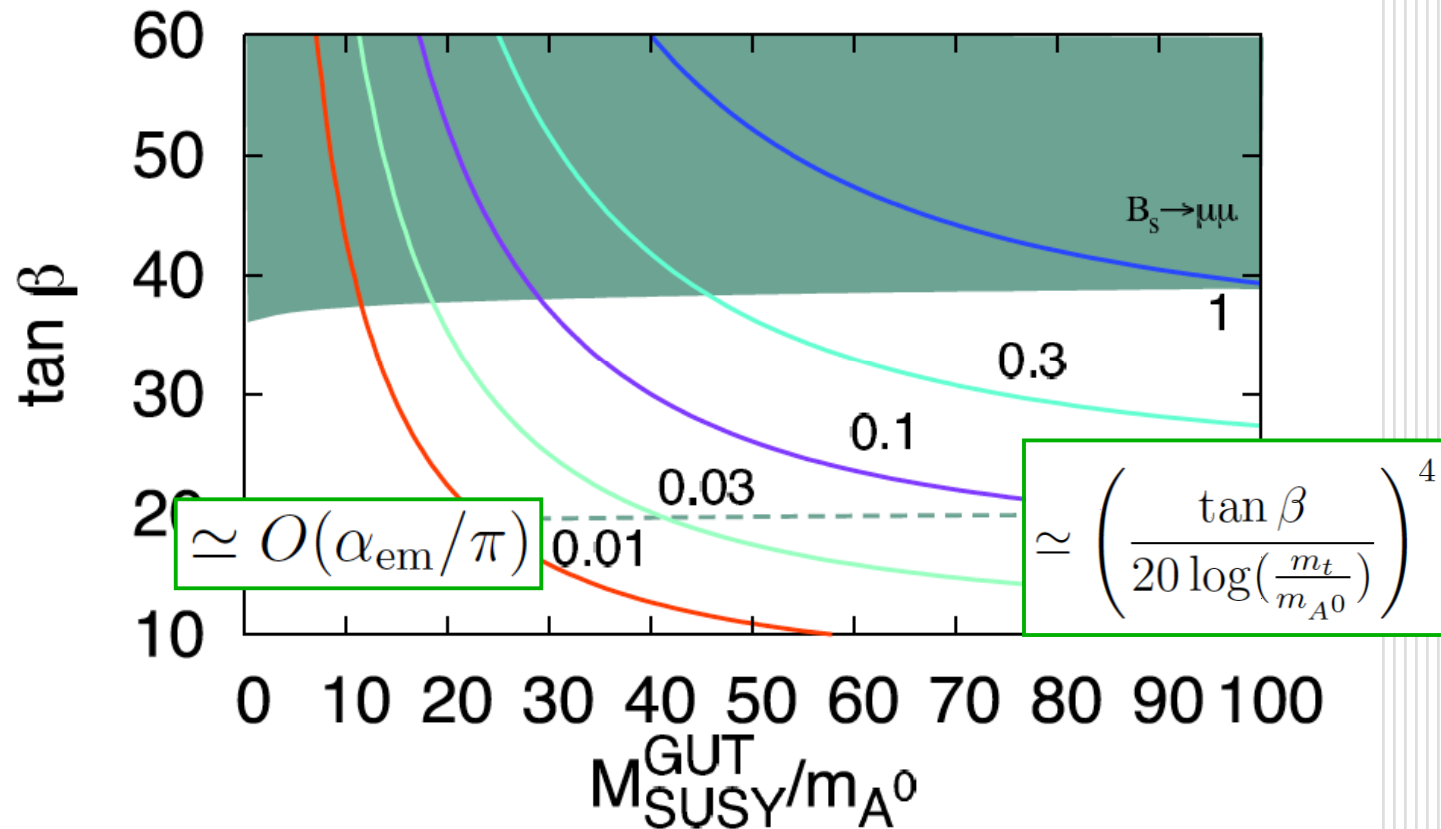


$$\simeq \left( \frac{\tan \beta}{20 \log\left(\frac{m_t}{m_{A^0}}\right)} \right)^4$$



# Left handed slepton mixing case

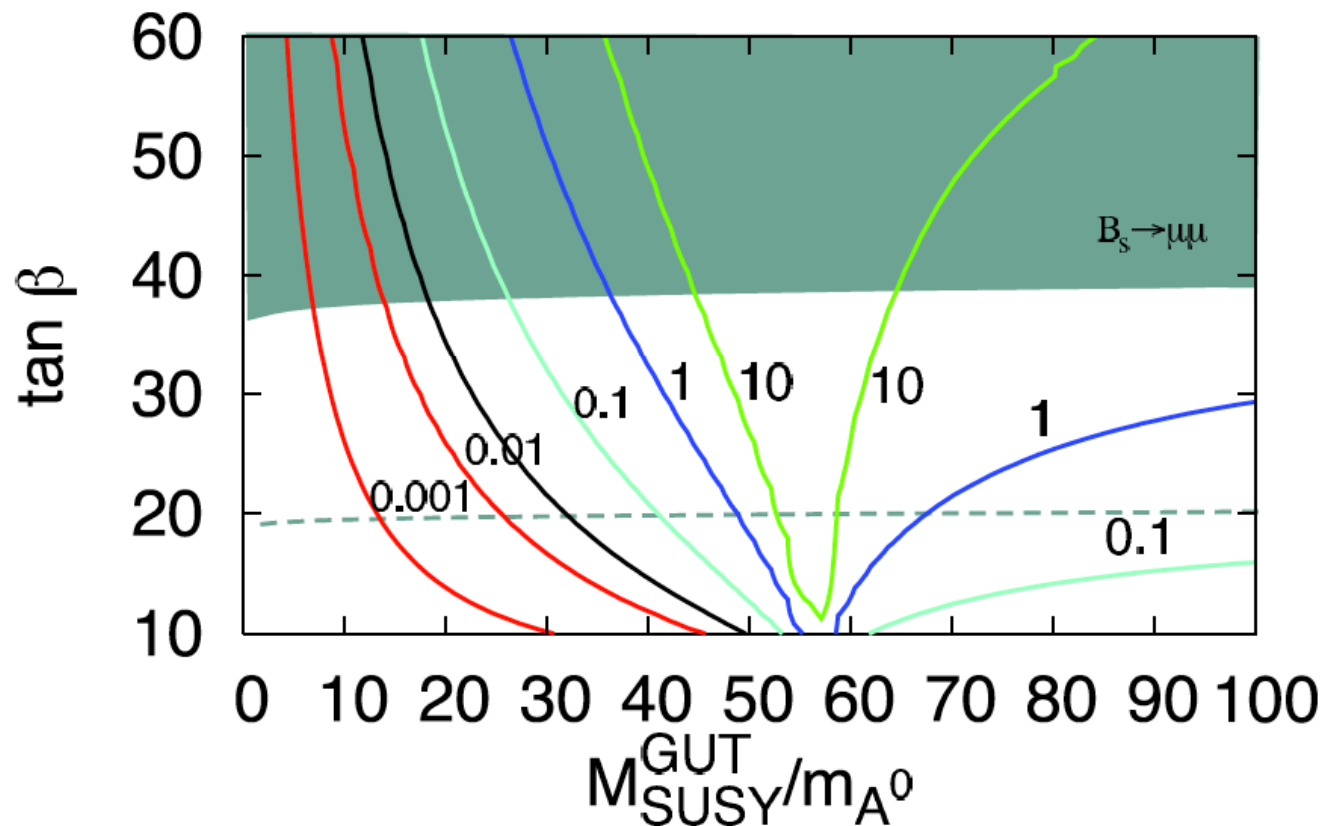
$$\text{BR}(\mu A^{\pm} \rightarrow e A^{\pm})^L / \text{BR}(\mu \rightarrow e \gamma)^L \quad \begin{array}{l} \mu = \mu_{\text{SUGRA}} \quad m_{A^0} = 500 [\text{GeV}] \\ A_0 = 0 [\text{GeV}] \quad m_0 = M_{1/2} = M_{\text{SUSY}}^{\text{GUT}} \end{array}$$



cLFV correlation gives information of **Higgs and SUSY sector**.

# Right handed slepton mixing case

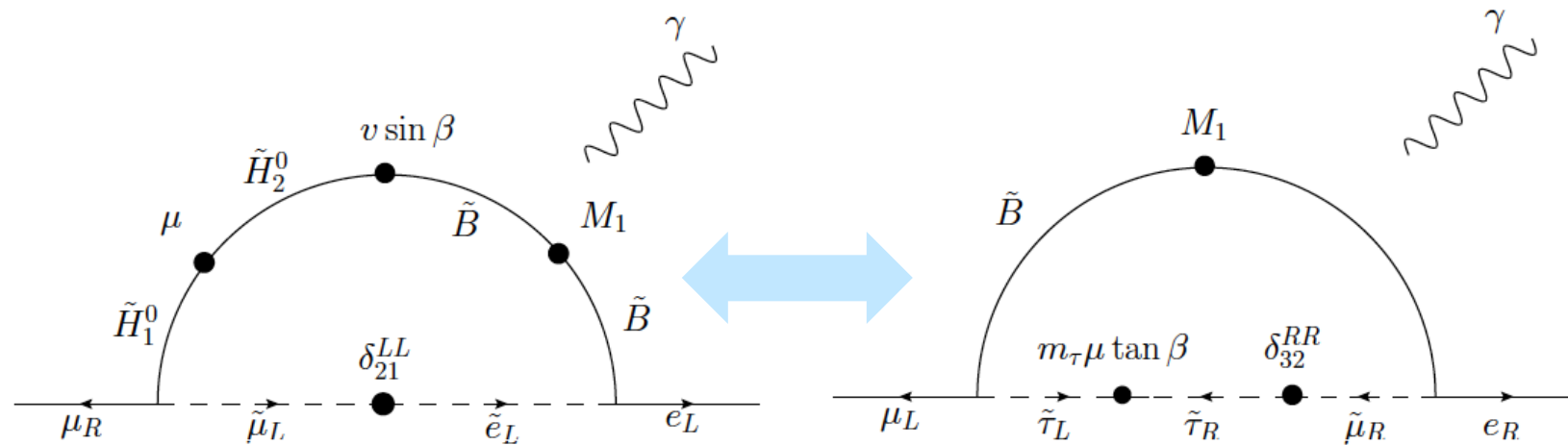
$$\text{BR}(\mu A \rightarrow e A) / \text{BR}(\mu \rightarrow e \gamma) \quad \mu = \mu_{\text{SUGRA}} \quad m_{A^0} = 500 [\text{GeV}] \\ A_0 = 0 [\text{GeV}] \quad m_0 = M_{1/2} = M_{\text{SUSY}}^{\text{GUT}}$$



Cancellation structure is different from left-handed mixing case.

# Right handed slepton mixing case

- When only  $\tilde{e}_{Ri}$  have flavor mixing, there are cancellation in the  $\mu \rightarrow e \gamma$  process.

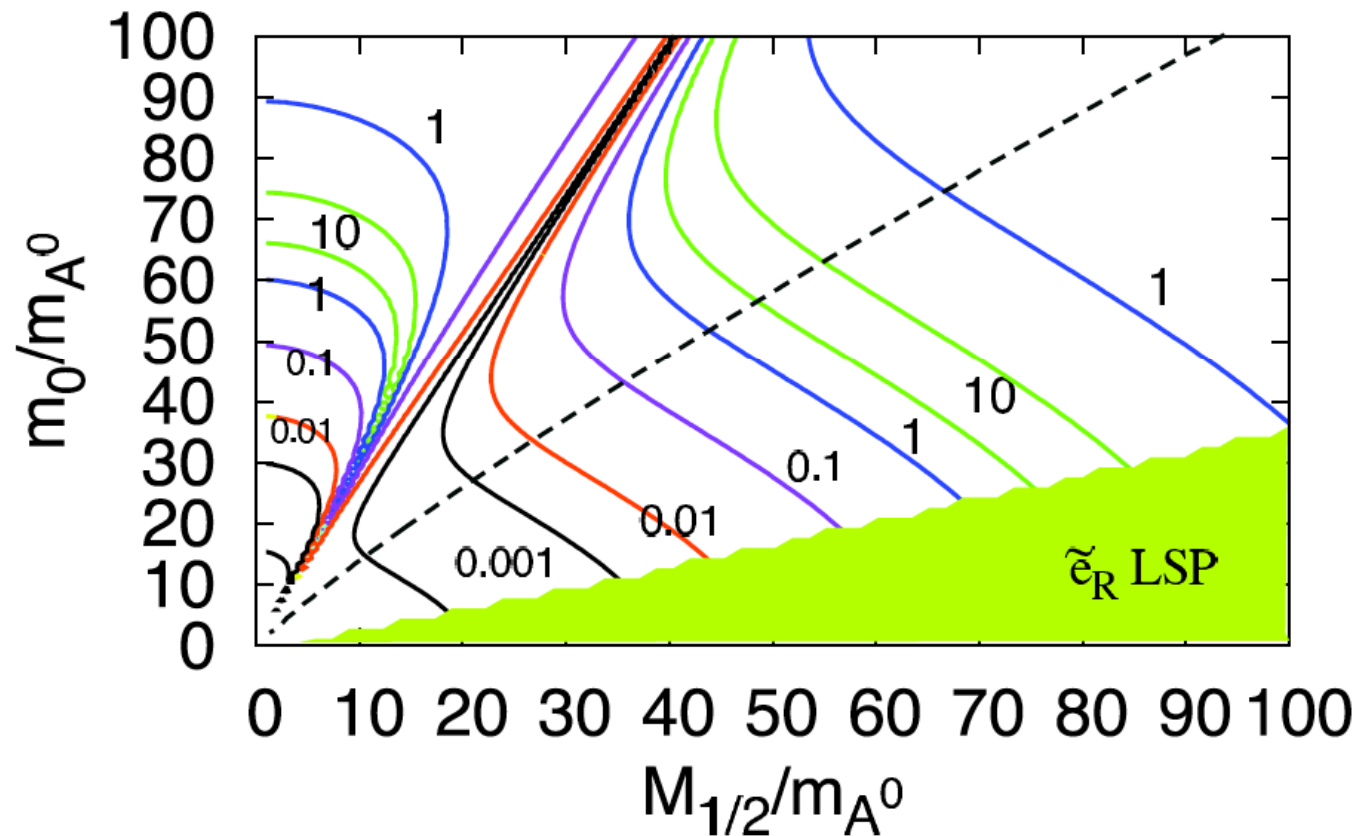


**Always destructive** (Hisano, Moroi, Tobe, Yamaguchi)  
and complete cancel when  $\mu \simeq \tilde{m}_L$

## Right handed slepton mixing case

- When only  $\tilde{e}_{Ri}$  have flavor mixing, there are cancellation in the  $\mu \rightarrow e\gamma$  process.

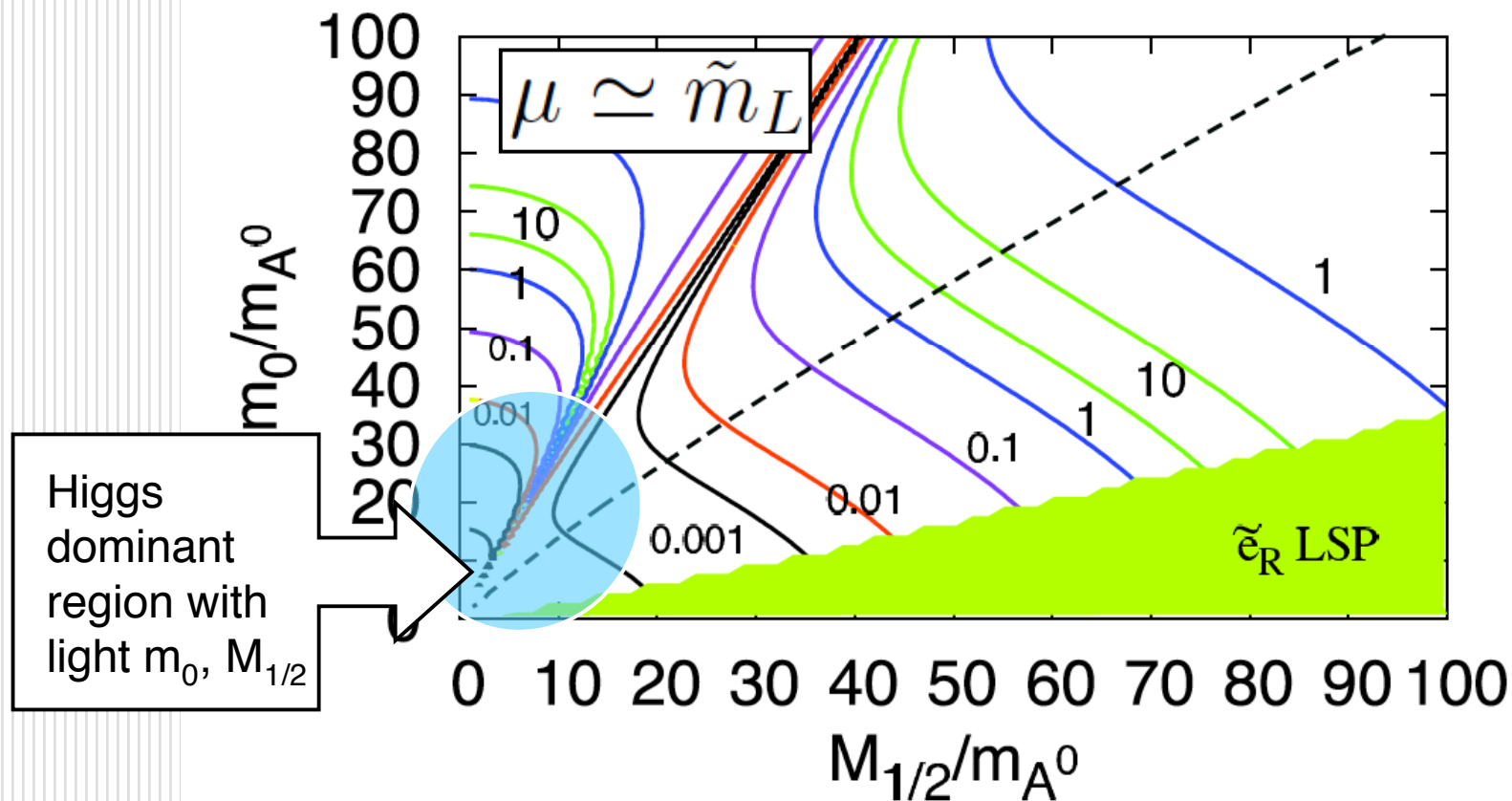
$$\text{BR}(\mu N \rightarrow e N)^R / \text{BR}(\mu \rightarrow e\gamma)^R \quad \begin{array}{l} \mu = \mu_{\text{SUGRA}} \quad A_0 = 0 [\text{GeV}] \\ m_{A^0} = 500 [\text{GeV}] \quad \tan\beta = 20 \end{array}$$



# Right handed slepton mixing case

- When only  $\tilde{e}_{Ri}$  have flavor mixing, there are cancellation in the  $\mu \rightarrow e\gamma$  process.

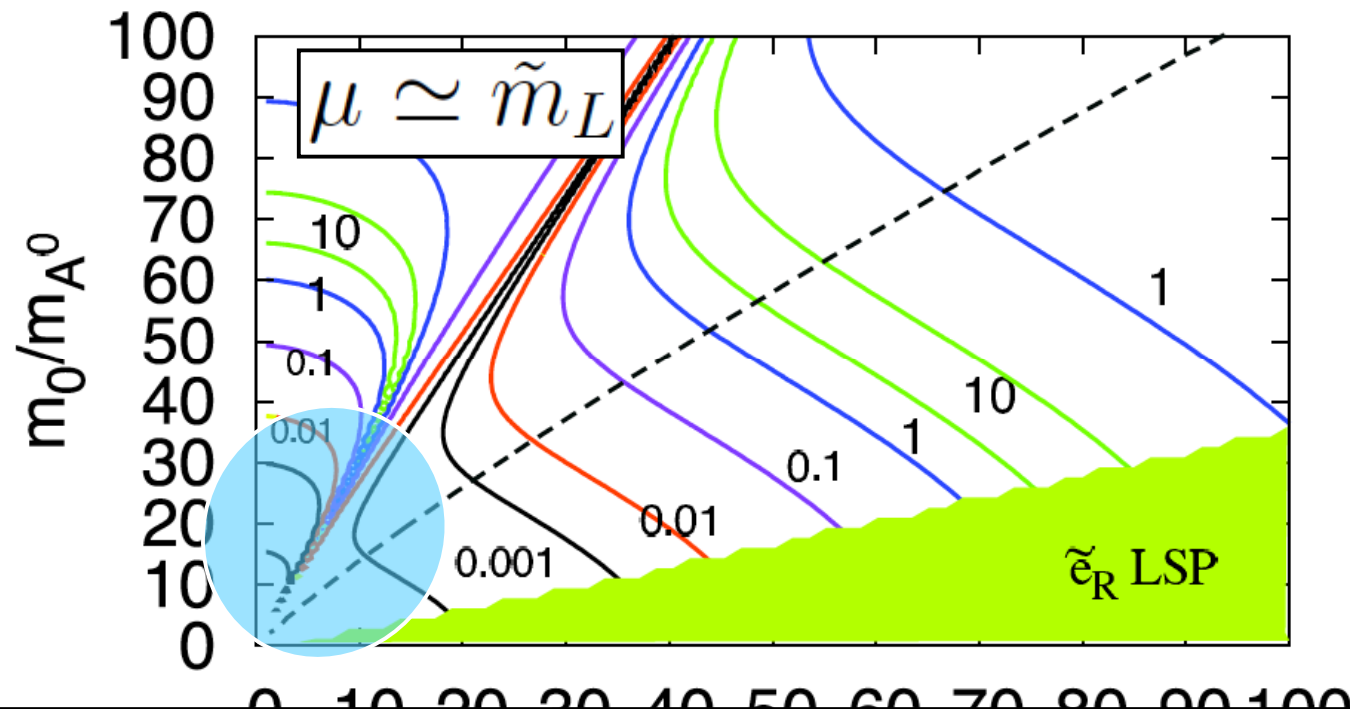
$$\text{BR}(\mu N \rightarrow e N)^R / \text{BR}(\mu \rightarrow e\gamma)^R \quad \begin{array}{l} \mu = \mu_{\text{SUGRA}} \quad A_0 = 0 [\text{GeV}] \\ m_{A^0} = 500 [\text{GeV}] \quad \tan\beta = 20 \end{array}$$



# Right handed slepton mixing case

- When only  $\tilde{e}_{Ri}$  have flavor mixing, there are cancellation in the  $\mu \rightarrow e\gamma$  process.

$$\text{BR}(\mu N \rightarrow e N)^R / \text{BR}(\mu \rightarrow e\gamma)^R \quad \begin{array}{l} \mu = \mu_{\text{SUGRA}} \quad A_0 = 0 [\text{GeV}] \\ m_{A^0} = 500 [\text{GeV}] \quad \tan\beta = 20 \end{array}$$



Although  $m_0$  and  $M_{1/2}$  are small, the ratio deviates sizably from  $O(\alpha)$ .

# Conclusion

- We study non-decoupling  $\mu$ -e transition effects by Higgs-mediated contribution in the MSSM.
- It is considered that the Higgs effect becomes dominant when SUSY particles are decouple. However, there are some **Higgs dominant region** although SUSY particle masses are TeV scale.
- $\text{BR}(\mu N \rightarrow e N) / \text{BR}(\mu \rightarrow e \gamma)$  could constrain  $\tan\beta$ ,  $m_A$ , or  $M_{\text{SUSY}}$ . The cLFV correlation gives information of **Higgs sector** or **SUSY sector**.





**That's all.  
Thank you.**